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VAPORIZATION OF COMPOUNDS AND ALLOYS AT HIGH TEMPERATURES

PART X. THE DISSOCIATION ENERGY OF THE GROUP IV — GROUP VI MOLECULES

TECHNICAL REPORT NO. WADD-TR-60-782, PART X May 1963

DIRECTORATE OF MATERIALS AND PROCESSES AERONAUTICAL SYSTEMS DIVISION AIR FORCE SYSTEMS COMMAND WRIGHT-PATTERSON AIR FORCE BASE, OHIO

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(Prepared under Contract No. AF61(052)-225 by the Université Libre de Bruxelles, Brussels, Belgium; R. Colin and J. Drowart)

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FOREWORD

This report was prepared by the University of Brussels, Belgium, under USAF Contract No. AF61(052)-225. The contract was initiated under Project No. 7350, "Refractory Inorganic Non-Metallic Materials," Task No. 735001, "Non-Graphitic." The work was administered under the direction of the Directorate of Materials and Processes, Deputy for Technology, Aeronautical Systems Division, Wright-Patterson Air Force Base, Ohio. Mr. F. W. Vahldiek was the project engineer.

ABSTRACT

The available spectroscopic and thermochemical data for the dissociation energy of the group IV-group VI MeX molecules are reviewed and the best present values proposed. Correlations of several excited molecular states with atomic products are proposed.

This technical documentary report has been reviewed and is approved.

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TABLE OF CONTENTS

THE	DISSOCIATION ENERGY OF THE GROUP IV-GROUP VI MOLECULES	:
	Thermochemical Data	2
	Spectroscopic Data	16
	Comparison of Thermochemical and Spectroscopic Data	21
	Acknowledgments	27
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THE DISSOCIATION ENERGY OF THE GROUP IV-GROUP VI MOLECULES

Accurate convergence limits for one or more electronic states of most of the group IV (Me: C, Si, Ge, Sn, Pb)group VI (X: 0, S, Se, Te) molecules are known. This is in particular the case for CO (1), a molecule not further discussed here and SiS, SiSe, GeS, SnO, SnS, SnSe and PbS, molecules for which the best known convergence limit is that of the excited E state. Accurate spectroscopic values $D_{0}^{"}$ for the dissociation energy of the ground state of these molecules can be derived therefrom, provided the correlation with atomic products is known. It has been postulated (2,3) that the E states of these molecules are analogous and correlate with the same atomic products. Recent rotational analysis for PbO⁽⁴⁾ and SnO⁽⁵⁾ showed that for the E state of these molecules, the most likely correlation is $Me(^{3}P_{1})+0(^{3}P_{1})$. Likewise the comparison of thermochemical and spectroscopic data for SnS and PbS (6) showed that their E state most probably correlates with $Me(^{3}P_{1}) + S(^{3}P_{1...0})$.

The purpose of this paper is to present a comparison between the spectroscopic data and recent thermochemical determinations of the dissociation energies of the group IV-group Manuscript released by the authors February 1963 for publication as an ASD Documentary Report.

VI molecules, several of which were performed mass spectrometrically and to derive in each case the best present value for the dissociation energy. The comparison gives further support to the correlation of the E state with $Me(^3P_1) + X(^3P_1)$, which is then applied to all molecules of this group for which reliable data for this state are known. Other available spectroscopic data are also examined.

THERMOCHEMICAL DATA

The thermochemical data were in general obtained mass spectrometrically or from total pressure measurements supplemented with mass spectrometric analysis of the composition of the vapor. The mass spectrometric technique (7) and the instruments (8-10) used have been described previously.

The system studied, the temperature intervals covered and the composition of the vapor are summarized in Table I. This table shows that above MeX compounds, including SiO, GeO and SnO known to disproportionate to 1/2 MeO₂(s)+1/2Me(s,1), the predominent molecule is MeX, polymers (and elements for GeTe, SnSe, SnTe, PbS, PbSe and PbTe) representing in general 10⁻¹ to 10⁻⁴ in pressure, except for SnO and PbO, where (SnO)₁₋₄ and (PbO)₁₋₄ have comparable pressures. The MeX₂ compounds vaporize mainly to MeX(g)+1/2X₂(g). MeX₂ molecules were observed above SiO₂⁽¹¹⁾, SiS₂, GeTe and SnTe.

TABLE I: Gaseous Species in Equilibrium with Group IV-Group VI Compouds Standard Heats of Formation and of Sublimation

1200 - 1460 11 108.9*0.1 13	System	Composition of the vapour	Temperature range investi- gated (°K)	Ref.	$-\Delta H_{298,f}^{\bullet}$ kcal/mole	Ref.	AH898, subl(b)	(b) Ref.
Si0, 0, Si02, 02 1750 - 1950 11 217.7 #0.2 13 SiS 850 - 1140 38.7 #3.0 15 1140 15.8 #3.0 13 SiS, S2, SiS2 850 - 1140 61.5 #3.0 61.5 #3.0 10 102 + 1140 61.5 #3.0 14 GeO, GeO)2, GeO)3, GeO)3, GeO)3, GeO)3, GeO)3, GeO)3, GeO 1020 - 1520 128.3 #0.5 14 GeS	SiO ₂ +1/2	SiO,	i	11	108.9*0.1		. 2	20
sis 850 - 1140 38.7*3.0 sis2 sis2, sis2 850 - 1140 61.5*3.0 GeO (GeO,(GeO)2,(GeO)3) - 58.9*2.5 14 GeO2 GeO,(GeO)2,(GeO)3 750 - 950 128.3*0.5 14 GeO2 GeO, O2 1020 - 1520 128.3*0.5 14 GeS - - - 16.2*1.3 14 GeS - - - 16.2*1.3 14 GeS - - - 16.2*1.3 14 GeS -	\sin_2	0, SiO ₂ ,	1 19	11	217.7±0.2	13		21
SiS2 SiS, S2, SiS2 850 - 1140 61.5*3.0 GeO (GeO, (GeO)2, (GeO)3) - 58.9\$2.5 1/2GeO2+1/2Ge GeO, (GeO)2, (GeO)3 750 - 950 64.2\$*0.3 14 GeO2 GeO, O2 1020 - 1520 128.3\$*0.5 14 GeS - - 1020 - 1520 16.2\$*1.3 14 GeS - - 1000 - 1520 16.2\$*1.3 14 GeS - - 1000 - 1500 16.2\$*1.3 15 SnO2 SnO, (SnO)2-4 1100 - 1400 138.7\$*0.2 16 SnS SnS, (SnS)2 835 - 1005 6 25.1\$*1.2 17 SnS SnS, (SnS)2 790 - 1010 19.9\$*2.0 18 PbO PbO, (PbO)2-4 1000 - 1800 14.6\$*0.3 18 PbO PbO, (PbO)2-4 1000 - 1500 19.9\$*1.0 19.9\$*2.0 19 SnS SnS, (SnS)2 835 - 1005 6 25.1\$*1.0 19 19.9\$*2.0 19 PbO PbO, (PbO)2-4 830 - 990 14.6\$*6.0 19 19 19 1	Sis		50 -		8.7±3.		#	
(GeO, (GeO) ₂ , (GeO) ₃) - 58.9±2.5 GeO, GeO, CGeO) ₂ , (GeO) ₃ GeO, O ₂ GeO, O ₂ GeC, O ₃	sis ₂	S ₂ ,	50 -		1.5±3.		5	
1/2 Geo ₂ +1/2 Ge Geo, (Geo) ₂ , (Geo) ₃ 750 - 950 64,2 ± 0.3 14 Geo ₂ Geo, O ₂ 1020 - 1520 128,3 ± 0.5 14 GeS - 152 - 152 16,2 ± 1.3 15 GeS - 102 - 150 16,2 ± 1.3 15 GeTe GeTe, Te ₂ , GeTe ₂ 870 - 1400 4,2 ± 2.0 16 1/2 Sno ₂ +1/2 Sn Sno, (Sno) ₂ -4 1100 - 1400 188.7 ± 0.2 16 SnS Sno, O ₂ , (Sno) ₂ 1200 - 1500 138.7 ± 0.2 16 SnS SnS, (SnS) ₂ 835 - 1005 6 25.1 ± 1.2 17 SnTe SnSe, (SnSe) ₂ , SnTe ₂ 790 - 1010 14,6 ± 0.3 18 PbO PbO, (PbO) ₂ -4 1800 - 1180 52.4 ± 0.2 19 PbS PbS, (PbS) ₂ , Pb, S ₂ 1865 - 1140 6 22.5 ± 0.5 19 PbS PbSe, Pb, Se ₂ 1180 - 1450 12 18.0 ± 2.0 19	GeO	(GeO, (GeO) ₂ , (GeO) ₃)	ı		8.9*2.		50,6*2,5	
GeO, O ₂ GeS - 1520 GeTe, Te ₂ ,GeTe ₂ Solo, O ₂ (SnO) ₂ -4 SnO, O ₂ (SnO) ₂ SnS, (SnS) ₂ SnSe, (SnSe) ₂ SnSe, (1/2GeO ₂ +1/2Ge	GeO,(GeO) ₂ ,(GeO) ₃	o		•	1,4	5	
GetS GetS - - 16.2*1.3 15 38.7*0 GeTe GeTe, Te ₂ ,GeTe ₂ 870 - 954 4,2*2.0 16 46.6*2 1/2Sn0 ₂ +1/2Sn Sn0,(Sn0) ₂ -4 1100 - 1400 99.8*0.1 16 71.1*1 Sn0 ₂ Sn0, O ₂ ,(Sn0) ₂ 835 - 1005 6 25.1*1.2 17 52.6*1 SnS SnSe,(SnS) ₂ 835 - 1005 6 25.1*1.2 17 52.6*1 SnSe SnSe,(SnS) ₂ 830 - 1010 19.9*2.0 1 52.6*1 PbO PbO,(PbO) ₂ -4 830 - 1010 14.6*0.3 18 52.5*1 PbS PbS,(PbS) ₂ ,Pb,S ₂ 940 - 1180 52.4*0.2 19 66.5*2 PbS PbS,(PbS) ₂ ,Pb,S ₂ 865 - 1140 6 22.5*0.5 19 55.7*1 PbSe PbSe,Pb,Se ₂ 1180 - 1450 12 18.0*2.0 18 53.9*1 PbTe PbTe,Pb,Te,Te,Te 923 - 1171 12 16.6*0.5 18 53.8*2	GeO ₂		-		•	7 ;	(120,0*2,0)	
GeTe, Te ₂ , GeTe ₂ 870 - 954 $_{02}+1/2Sn$ Sn0, $_{02}, (Sn0)_{2-4}$ 1100 - 1400 Sn0, $_{02}, (Sn0)_{2-4}$ Sn0, $_{02}, (Sn0)_{2}$ SnS, $_{(SnS)_{2}}$ SnS, $_{(SnS)_{2}}$ SnS, $_{(SnS)_{2}}$ SnS, $_{(SnS)_{2}}$ SnSe, $_{(SnSe)_{2}}$, Se ₂ SnTe, Te, Te, Te, Te, SnTe ₂ SnTe, Te, Te, Te, SnTe ₂ SnSe, $_{(SnSe)_{2-4}}$ PbO, $_{(PbO)_{2-4}}$ PbS, $_{(PbS)_{2}}$ PbS, $_{(PbS)_{2}}$ PbS, $_{(PbS)_{2}}$ PbS, $_{(PbS)_{2}}$ PbS, $_{(PbS)_{2}}$ PbSe, Pb, Se ₂ PbS, $_{(PbS)_{2}}$ PbSe, Pb, Se ₂ PbSe, PbSe, PbSe, PbSe, Se ₂ PbSe,		GeS	!		.2*1.	15	38,7*0.6	22
02+1/2Sn Sn0,(Sn0) ₂ -4 1100 - 1400 69.8*0.1 16 71.1*1 Sn0, 0 ₂ ,(Sn0) ₂ SnS,(SnS) ₂ SnSe,(SnSe) ₂ , Se ₂ SnTe, Te, Te, Te ₂ , SnTe ₂ , SnTe ₂ Pb0,(Pb0) ₂ -4 PbS,(PbS) ₂ ,Pb,S ₂ PbSe,Pb,Se ₂ SnO - 1010 19.9*2.0 14.6*0.3 18 52.6*1 51.1*2 52.6*1 52.6*1 52.6*1 52.6*1 52.6*1 52.6*1 52.6*1 52.6*1 FbS,(PbS) ₂ ,Pb,S ₂ FbS,(PbS) ₂ ,Pb,S ₂ 1180 - 1450 FbSe,Pb,Se ₂ 1180 - 1450 12 18.0*2.0 18 53.9*1	GeTe	GeTe, Te $_2$, GeTe $_2$	ග 		.2±2.	91	ę•e	23,24
SnO, O ₂ ,(SnO) ₂ SnS,(SnS) ₂ SnSe,(SnS) ₂ SnSe,(SnSe) ₂ , Se ₂ SnTe, Te, Te ₃ , Te ₂ SnTe, Te, Te ₃ , SnTe ₂ SnTe, Te, Te ₃ , SnTe ₂ SnTe, Te, Te ₃ , SnTe ₂ SnTe, Te, Te ₃ SnSe,(SnSe) ₂ , Se ₂ SnTe, Te, Te ₃ SnSe,(SnSe) ₂ SnSe,(S	1/2SnO ₂ +1/2Sn	SnO,(SnO) ₂₋₄	ı		0 * 8 * 0	16	71,1*1,3	
SnS,(SnS) ₂ SnSe,(SnSe) ₂ , Se ₂ Tg, (SnTe) ₂ SnTe, Te, Te, Te ₂ SnTe, Te, Te ₂ SnTe, Te, Te ₂ SnTe, Te, Te ₂ SnTe, Te, Te ₃ Pb0,(Pb0) ₂₋₄ PbS,(PbS) ₂ ,Pb,S ₂ PbSe,Pb,Se ₂ Tille - 1450 Till	sno_2		ŧ		•	16	139,7*2.0	
SnSe,(SnSe) ₂ , Se ₂ 790 - 1010 19.9±2.0 51.1±2 SnTe, Te, Te ₂ , SnTe ₂ , 830 - 990 14.6±0.3 18 52.5±1 (SnTe) ₂ PbO,(PbO) ₂₋₄ 940 - 1180 52.4±0.2 19 66.5±2 PbS,(PbS) ₂ ,Pb,S ₂ 865 - 1140 6 22.5±0.5 19 55.7±1 PbSe,Pb,Se ₂ 1180 - 1450 12 18.0±2.0 18 53.9±1 PbTe,Pb,Te,Te ₂ 923 - 1171 12 16.6±0.5 18 53.8±2	SnS	SnS,(SnS) ₂	35 -	9	5,1*1.	17	52.6*1.6	ဖ
SnTe, Te, Te, Te, SnTe ₂ , SnTe ₃ , SnT	SnSe		1		• თ		51.1 \$ 2.0	a,25
PbO,(PbO), (PbO), (PbO), PbO, PbO, PbO, PbO, PbO, PbO, PbO, PbO	SnTe	Te, Te,	830 - 8		•	18	2.5±	a,25
PbS,(PbS) ₂ ,Pb,S ₂ 8 66 - 1140 6 22.5±0.5 19 55.7±1 PbSe,Pb,Se ₂ 1180 - 1450 12 18.0±2.0 18 53.9±1 PbTe,Pb,Te,Te,	Pbo	Pb0, (Pb0),	ı		52.4±0.2	19	66.5*2.5	
PbSe,Pb,Se ₂ 1180 - 1450 12 18.0*2.0 18 53.9*1 PbTe,Pb,Te,Te,	PbS	PbS, (PbS)2, Pb, S2	ı	9	2.5±0.	19	55.7*1,6	9
PbTe, Pb, Te, Te, Te, 923 - 1171 12 16.6*0.5 18 53.8*2	PbSe	PbSe, Pb, Se ₂	- 1450	12	18.0*2.0	18	3.9±1	26
	PbTe	Pbre, Pb, re, re,	၊ ဗ	12	16.6±0.5	18	53.8 ± 2.0	27

Legend of Table I:

- a. this work
- b. ΔH° refers to the equilibrium 298, subl

 $MeX(s) \rightarrow MeX(g)$ or

 $MeX_2(s) \rightarrow MeX(g)+1/2X_2(g)$

The dissociation energies $D_{o}^{O}(\text{MeX})$ were calculated from thermochemical cycles of the type

Additional determinations based on the equilibrium $MeX(g) + Me(s, 1 \text{ or } g) + 1/2 X_2(g)$ were made for GeTe, SnSe SnTe, PbS⁽⁶⁾, PbSe⁽¹²⁾ and PbTe⁽¹²⁾.

The heats of sublimation $\Delta H_{\rm sub}({\rm MeX})$ are summarized in Table II. The heats of sublimation of the group IV elements Me⁽²⁸⁾ and of the group VI molecules $X_2^{(28)}$ as well as the heats of formation of the compounds MeX or MeX₂ were in general taken from the litterature. The values used for the dissociation energies $D_{298}^{\rm O}$ of $D_{298}^{\rm O}$, $D_{298}^{$

(It may be noted that the agreement between spectroscopic and thermochemical values of the dissociation energies of the group of molecules discussed in the following chapters is an indirect confirmation of the $D_O^O(S_2)$, $D_O^O(S_2)$ and $D_O^O(Te_2)$ values used).

For all the gaseous species Me, X2 and MeX, values for free energy function given in the literature (28,32) were used. For those stable compounds MeX and MeX, for which values are not given in the literature (32,33) estimates were made as follows. For solid SnTe, PbSe and PbTe (32), the entropies at 298°K are known. For these compounds only the heat capacity, while for SiS, GeTe and SnSe both the entropy at 298°K and the heat capacity were estimated from the corresponding quantities for the elements using approximations discussed by Kubaschewski and Evans (18). The corresponding quantities for GeS were derived from thermodynamic data for the reactions GeS(s) \rightarrow GeS(g) (22) and GeS(s) \rightarrow Ge(s) + 1/2 S₂ (15). The values obtained for the free energy functions are at 298°K, 12.0(SiS), 18.6(GeS), 20.3(GeSe), 22.2(GeTe), 22.5 (SnSe) and at 900°K 16.2(SiS), 23.0(GeS), 24.8(GeSe), 26.8 (GeTe), 28.0(SnSe), 29.8(SnTe), 30.0(PbSe) and 31.8(PbTe) cal/degree/mole.

The thermochemical data for each molecule are briefly discussed below

CS Thermochemical determinations (34) of the heat of formation of CS(g) were based on the reactions

C(graphite) +
$$CS_2(g)$$
 + 2 $CS(g)$
MnS(s) + C(graphite) + $CS(g)$ + $Mn(g)$

Using the latter reaction for which the enthalpy change ΔH_{298}^{o} was measured to be 178.5±0.5 kcal/mole (34), $\Delta H_{298,f}^{o}$ (MnS(s) = 49.5±0.5 kcal/mole and ΔH_{298}^{o} (sub Mn) = 66.6±1.0 kcal/

mole $^{(28,30)}$, Freeman $^{(37)}$ calculated a value of 62.5 ± 3.0 kcal/mole for the standard heat of formation of CS(g). This leads to a value $D_o^o(CS)=173.6\pm 3.5$ kcal/mole $^{(7.53\pm 0.15 \text{ eV})}$.

- CSe Thermochemical data were obtained (38) for the reaction C(graphite)+Se(g) \rightarrow CSe(g) Δ H°=-31.4±2.0 kcal/mole whence D°(CSe)=138.4±2.5 kcal/mole (6.00±0.10 eV)
- CTe This molecule has hiterto not been observed. Its dissociation energy is estimated here as D₀°=111±9 kcal/mole (4.8±0.4 eV) by comparison with the other IV-VI molecules.
- Sio The mass spectrometric data (11) confirm earlier conclusions that the main vaporization processes are

$$1/2 \operatorname{Si0}_{2}(s) + 1/2 \operatorname{Si}(s,1) + \operatorname{Si0}(g)$$
 I
 $\operatorname{Si0}_{2}(s) + \operatorname{Si0}(g) + 1/2 \operatorname{O}_{2}(g)$ II

The heats of these reactions were taken as $\Delta H_{298}=85.5$ $\pm 0.5^{(20)}$ and $\Delta H_{298}^{\circ}=198.3\pm 1.5^{(21)}$ kcal/mole respectively. To derive the heat of formation of SiO(g), the reaction $^{(20)}$ SiO₂(s)+H₂(g) + SiO(g)+H₂O(g)(III), $\Delta H_{298}^{\circ}=134.8\pm 2.0$ kcal/mole was further taken into account. (For a more complete review, see Kubaschewski and Evans and Schick $^{(35)}$). The determination for reaction I and III are in good agreement with one another, while those for reaction II lead to slightly higher values. The recently redetermined standard heat of formation of SiO₂, ΔH_{298f}° (SiO₂(s))=-217.7±0.2 kcal/mole $^{(13)}$ was used

and the heat of sublimation of silicon taken as $\Delta H_{298,sub}^{0}(Si) = 107.9 \pm 1.0$, this being the average of several determinations (40). The resulting mean value for the dissociation energy of SiO(g) is $D_{0}^{0}(SiO) = 191.9 \pm 3.0$ kcal/mole (8.32 ± 0.13 eV).

SiS No independant thermochemical data for this molecule were yet obtained. The spectroscopic value for the dissociation energy $D_0'' = 6.38\pm0.06$ eV was used to calculate the heats of formation of SiS(s) and SiS₂(s) from mass spectrometric study of mixtures of SiS and SiS₂ and from total pressure for SiS₂(41). The heats of formation obtained in this manner from the reactions

SiS(s) + SiS(g) $\Delta H_{298}^{\circ} = 63.4 \pm 3 \text{ kcal/mole}$ SiS₂(s) + SiS(g)+1/2S₂(g) $\Delta H_{298}^{\circ} = 101.5 \pm 3 \text{ kcal/mole}$ are ΔH_{298f}° (SiS(s)) = -38.8 \pm 3.0, and ΔH_{298f}° (SiS₂(s))= -61.9 \pm 3.0 kcal/mole

The heat of formation of condensed SiS shows this compound not to be metastable up to about 1250° K, in agreement with observations by Retzlaff and Kohlmeyer (42). The heat of formation of SiS_2 is quite different from the discordant values given in the literature, which have anyway to be corrected for the recent values for the heat of formation of SiO_2 . The average value

selected by Kubaschewski and Evans $^{(18)}$, 61.6±6.0 kcal/mole (corrected for $\Delta H_f(SiO_2)$) is in agreement with the present value. The latter also makes it possible to explain the apparent boiling points of SiS_2 at about 1400°K as being that of the $Si+SiS_2$ system.

GeO Earlier thermochemical values were based on vapor pressure determination for GeO (43) and Ge + GeO (14,43)

The evaporation behavior of Ge + GeO and GeO were examined mass spectrometrically. The evaporation processes are

 $1/2\text{Ge}(s) + 1/2\text{GeO}(s) + 1/\text{m}(\text{GeO})_{n}(g) \quad n=1,2,3$ $\text{GeO}_{2}(s) + \text{GeO}(g) + 1/2\text{O}_{2}(g)$

The mass spectrometric results for GeO_2 obtained here are in disagreement with an earlier investigation (44), also by mass spectrometry, which seems to be characteristic of $Ge + GeO_2$ rather than of GeO_2 . The enthalpy for the reaction

 $1/2 \text{GeO}_2(s) + 1/2 \text{Ge}(s) + \text{GeO}(g) \Delta \text{H}_{298}^{\circ} = 55.9 \pm 1.6 \text{ kcal/mole}$ was calculated from mass spectrometric pressure measurements made independently, as well as from the total pressures given in the literature $(1^4, 4^3)$ correcting these for the presence of $(\text{GeO})_2$ and $(\text{GeO})_3$. A similar procedure was adopted for GeO(s), for which a heat of formation $\Delta \text{H}_{298}^{\circ} = -58.9 \pm 3.0 \text{ kcal/mole}$ was

derived from the earlier data (43) in satisfactory agreement with direct determinations (45). The best average of all data gives $D_0^0(\text{GeO}) = 156.6 \pm 3.0 \text{ kcal/}$ mole (6.79 \pm 0.13 eV). The enthalpy of the reaction $\text{GeO}_2(s) + \text{GeO}(g) + 1/2 \cdot 0_2(g)$ was not measured. As a result of interaction with SiO_2 crucible the pressures in the literature (46) are low. The enthalpy given in Table I was calculated from the heat of formation of $\text{GeO}_2(s)$ and the above dissociation energy of GeO(g).

GeS The mass spectrometric analysis of the vapor showed only the monomer GeS in the gas phase.

The enthalpy of the reaction

GeS(s) + GeS(g) ΔH_{298}° = 38.7±0.6 kcal/mole is hence very well known from a number of pressures measurements which are all in very good agreement except that of Kenwonthy and coworkers which seems to be a factor of approximately 3 too high. The heat of formation of GeS was evaluated from the reaction (15)

2 GeS(s)
$$+$$
 2 Ge(s) $+$ S₂(g)

and yields

 ΔH_{298f}° (GeS) = -16.2±1.3 kcal/mole 298f Inserting these values in the cycle given above, one obtains D_{0}° (GeS) = 133.3±3.5 kcal/mole (5.78±0.15 eV).

- The total vapor pressure above GeSe was measured by
 Liu Ch'Un-Hua, Pashinkin and Novoselova (24)

 the spectroscopic value of the dissociation energy

 D"(GeSe) = 113.1±0.7 kcal/mole discussed below, it

 can be shown that the decomposition pressure of

 Se₂ is low compared to that of the main species GeSe(g).

 The total pressure measurements yield therefore a

 value of the heat of sublimation AH^o (GeSe) =

 298sub

 44.6±1.5 kcal/mole obtained by Ind (44.2) and 3rd

 law (44.8) procedures. Combining this value with

 D"(GeSe), a heat of formation AH^o (GeSe(s)) =

 14.4±3.0 kcal/mole is obtained.
 - The mass spectrometric investigation of the decomposition of the vapor showed the presence of GeTe,

 Te₂, (GeTe)₂ and GeTe₂. The pressure of the Te₂ molecule, relative to that of the main species GeTe(g) is 10⁻¹ at 880°K. Using total vapor pressure measurements (23,24) to evaluate the partial pressures, a heat of sublimation ΔH°₂₉₈, sub (GeTe) = 46.6±2.0 kcal/mole and a heat of formation ΔH°_{298f} (GeTe(s)) = -4.2±2.0 kcal/mole were calculated from 3rd law procedures (2nd law: ΔH°_{298sub} (GeTe) = 47.8 (23), 46.7 (24)). These values lead to the dissociation energy D°₀(GeTe)=92.4±3.0 kcal/mole (4.00±0.13 V.)

Sn0 The heats of both reactions

 $SnO_2(s) + SnO(g) + 1/2O_2(g)$

and $1/2\text{Sn}(1) + 1/2\text{SnO}_2(s) + 1/\text{n}(\text{SnO})_{n}(g)$ n=1,2,3,4 were studied mass spectrometrically. The values obtained are $\Delta H_{298}^{0} = 139.7 \pm 2.0$ and 71.1 ± 1.3 kcal/mole (for n=1) respectively. Whence $D_{0}^{0}(\text{SnO}) = 127.2 \pm 2.0$ kcal/mole (5.53 \pm 0.10 eV).

Total pressure determinations were examined for comparison. The vapor pressure measurements of $\mathrm{SnO}_2(s)$ by the Knudsen method (47) were not included, since the sample interacted with the crucible material (16) SiO_2 . The pressure determination by the transport method (16) gives too high values because the polymers $(\mathrm{SnO})_{2-4}$ contribute to the weight loss. Estimating their contribution from the mass spectrometer data yields $\mathrm{D}_0^\circ=130.5\pm3.0$ kcal/mole (5.64±0.13 eV) in agreement with the direct determinations.

SnS The thermochemical value for the dissociation energy based on mass spectrometry and vapor pressure data is $D_0^0(SnS) = 110.1\pm3.0 \text{ kcal/mole (4.78\pm0.13 eV)}$. For a detailed discussion see reference (6).

- Since The composition of the vapor as well as partial pressures were determined mass spectrometrically. The pressure of the main component, SnSe, was found in good agreement with pressure determinations by the effusion method (25). These data lead to a value for the heat of sublimation: ΔH° (SnSe)=51.1*2.0 kcal/298sub mole (3rd law). From the partial Se pressure, a heat of formation ΔH° (SnSe) = -19.9*1.5 kcal/mole was derived compared to the estimated value -16.5*2 kcal/mole (18). These values lead to D°(SnSe)=94.8*4.0 kcal/mole (4.11*0.17eV).
- <u>SnTe</u> The composition of the vapor and partial pressures were determined as above. The SnTe pressure was again in good agreement with total pressure measurements by the effusion method (25) The Te₂ partial pressure gave a value of -14.7±2.0 for the heat of formation of SnTe(s), in good agreement with ΔH⁰_{298f} =-14.6±0.3 kcal/mole (18), thereby confirming the values of the free energy function of SnTe(s) used. The 2nd and 3rd law values, 47.5 and 52.3 kcal/mole respectively, for the heat of sublimation are in less good agreement. Preference is given to the 3rd law value, which leads to D⁰(SnTe) = 79.8±3.0 kcal/mole (3.46±0.17 eV).

PbO The mass spectrometric analysis of the vapor above
PbO(s) showed the presence of (PbO)₁₋₄. The enthalpy
change for the reaction

 $PbO(s) \rightarrow PbO(g)$

was calculated as being ΔH_{298} = 65.8±1.5 kcal/mole from total pressure measurement given in the literature (48) taking the composition of the vapor into account. This value combined with the heat of formation of PbO(s), ΔH_{298f} = -52.4±0.2 kcal/mole (19) leads to D_0° (PbO) = 92.2±3.0 kcal/mole (4.00±0.13 eV).

- Mass spectrometric (6) and vapor pressure data combined with $\Delta H_{298f} = 22.5 \pm 0.5$ kcal/mole (18) lead to $D_0^0(PbS) = 78.9 \pm 0.25$ kcal/mole (3.42 \pm 0.12 eV). For a detailed discussion see reference (6).
- The knowledge of the dissociation energy of this molecule is not satisfactory. The mass spectrometric investigation of the vapor by Porter (12) shows PbSe to be the main gaseous species. Hence one calculates from the published pressure measurements (24) a heat of sublimation of 53.9 kcal/mole (2nd law: 53.2), which combined with the estimated heat of formation $\Delta H_{298f}^{O}(PbSe) = -18.0 \pm 2.0 \text{ kcal/mole} (18) \text{ gives } D_{0}^{O}(PbSe) = 65.0 \text{ kcal/mole} (2.82 \text{ eV}). \text{ Direct measurements made} in double oven type Knudsen cells by Porter (12) give$

for the reaction PbSe(g) \Rightarrow Pb(g) + Se(g) values ranging from 60.0 to 64.5 kcal/mole, with an average of 62.9 kcal/mole. Simultaneous measurements for the reaction Pb(g) + Se₂(g) \Rightarrow PbSe(g) + Se(g) give however $\Delta H_{298}^0 = 2.7 \pm 1.0 \text{ kcal/mole}$, which combined with $D_{298}^0 (\text{Se}_2) = 75.4 \pm 2.0 \text{ kcal/mole}^{(30)}$ would give $D_{298}^0 (\text{PbSe}) = 72.7 \pm 3.0 \text{ kcal/mole}^{(30)}$ would give $D_{298}^0 (\text{PbSe}) = 72.7 \pm 3.0 \text{ kcal/mole}^{(30)}$ of the heat of formation of PbSe(s) is available, an average value $D_0^0 (\text{PbSe}) = 63 \pm 5 \text{ kcal/mole}^{(2.95 \pm 0.22 \text{ eV})}$ is selected.

PbTe The dissociation energy of PbTe was measured mass spectrometrically by Porter⁽¹²⁾, who obtained $D_0^0(PbTe) = 51.4 \pm 2.0 \text{ kcal/mole.}$ This figure is in agreement with the value $D_0^0(PbTe) = 55.0 \pm 2.0 \text{ kcal/mole}$ mole (2nd law: 52.9) that can be derived from the vapor pressure measurements⁽²⁵⁾ (main component PbTe) and the heat of formation of PbTe(s), $\Delta H_{298f}^0 = -16.6 \pm 0.5 \text{ kcal/mole}^{(19)}$. The average value $D_0^0(PbTe) = 53.1 \pm 3.0 \text{ kcal/mole}^{(2.30 \pm 0.17 \text{ eV})}$ is selected.

The thermochemical values D_0^0 for the dissociation energies of the group IV-group VI molecules have been assembled in Table V.

SPECTROSCOPIC DATA

The spectroscopic data considered in the present paper have been collected in table II and were all taken from the literature. Only the following comments seem necessary for the purposes of this paper. For the molecules SiS, GeS, GeSe, GeTe, SnO and SnSe graphical extrapolations for the E state could be made. For GeS, GeSe, SnO, SnS, SnTe and PbS, the extrapolation were confirmed or supplemented by direct observation of convergence limits (continua) for the same or other states. For GeTe, Pb0 and possibly PbSe the value given for the convergence limit of the E state corresponds to the wavelength down to which vibrational structure lines could be observed. For the molecule GeTe the value of the convergence limit is somewhat uncertain, since it was estimated by analogy with GeS and GeSe. It should however not be lower than 4.38 eV, since the vibrational levels are known up to that energy. For the molecule SiTe, for which the linear Birge Sponer limit for the E state was not calculated because of the presence of large anharmonic and cubic terms, the extrapolation was nevertheless made here to serve as a guide for obtaining a spectroscopic value for the dissociation energy, since no other data are available.

TABLE II. Spectroscopic Data

Molecule	State	^	3	3 S	γe	ω ² / 4×ω	Conve	Convergence Limit	imit	
Mex		cm_L	cm_1	cm-1	cm-1	ev	linear ev	graph. ev	observ. ev	Ref
SO	׫	38804.8	1285.1	6.5 10.3		7.88 3.46	7.88 8.27			ნ #
ტ გ	׫	0 35134.6	1035.9 840	4.88		6.82	6.82			50
Sio	хошц	0 42640.4 52579.9 64497.2	1241.44 852.71 675 1116.5	5.92 6.44 4.15 7.22		8 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	8.07 8.78 9.91 13.34			ო
Sis	хош	0 34910.1 41750.8	749.5 513.12 403.54	2.56 2.93 1.40	-0.0329	6.92 2.79 3.61	6.92 7.12 8.79	44.9)		51 52
o Si Si Si Si	×ОШ	0 32360.2 38370.3	580.0 399.8 308.8	1.78 1.93 1.95	-0.032	5.86 1.52	5.86 6.57 6.28	5.66		23
SiTe	жды	0 28590.4 33871	481.2 338.6 242	1.30	1,30 1,70 (3,68)(+0,13)	5.52 2.09 (0.50)	5.52 5.64 (4.70)			53
0 9 9	хошы	0 37599.5 49396.5 67385	985.5 650 504.3 809.3	4.23 4.21 4.8 5.66		7.02 3.12 1.65 3.59	7.01 7.77 7.76 11.95			ო
ge S	×ОН	0 32589.2 38752.0	575.8 375.0 310.3	1.80		5.70 2.89 2.08	5.70 6.96 6.89	5.79	5.79	tt S
	;									

zw = 1.15

1 .

TABLE II. (Continued)

Molecule	State	o o o	3	æx	yω	ω ² /4×ω	Conver	Convergence Limit	mit	9.6
		cm_1	cm_1	cm-1	cm-1	e ^	linear	graph. ev	observ. ev	Ker.
GeSe	×	0	408.7	1.36		3.80	3,80			
	Ω	30776.0	269.4	∞		S	۳.			
	ы	5367.	17	0		.	Φ.	5,23	5.25	
GeTe	×	0	2	0.25		တ	თ			5.4
	А	27699.3	221.0	0.89		1.71	5.15			
	ш	139	7	(1.2)		7	ထ	6 † • †	≯ 4•38	
SnO	×		822.4	3.73		5.62	5.62			က
	щ	5325)	75)	4			•	q. :		
	ДΙ	950	582.8	3.08		3.42	7.08	(2.04)		
	щ	613	08.	•		. 7	• 5	5.68	•	
	0 6	~51775 57491							6.40-6.44	*
18	ንተ	875	724	21		0.77	8.05			
SnS	×	0	85.	35		က	ω,			55
	Q	8180.	30	1.265		2.51	6.01			
	ы	2939	295.0	•	-0.012	≠ .	ŝ	5.06		26
	Д	2174	S S							
	ದ ಗ	558470-890	0						₹7.05	
SnSe	×	0	31.	. 7		9	9			57,58
	٩ı	929	22	0.88		1.76	4.15			·
	m C	179269	2 2							
	DГ	27496.6	225.1	0.69	91000	2.28	5.69 3.69	i i		
	4 [14	7830)	900	•	9	•	•	† • †		

TABLE II. (Continued)

	Molecule State	00	3	:		3 4 - 3	linear	praph.	observ.	
	Cm	1	Cm T	cm -L	- ⊥ cm	o >	e <	0 0 0		
SnTe	×	0	259.5	0.50		ı —				57,58
	16	03.	178.5	†† . 0		2	3	,4		
	20	80.	230.3	1.53	-0.013	\circ	9	(3.34)		
	21	97.	218.1	86.0		S	Н			
	25	40	179.1	0,40		_	9			
	27	80.	135.0	2.5		2	9		(4,07)	
	G 284	65.0	0.86	1.0		0.30	3.83			
	29	42.	2000.8	0.3		~~1	α			
	30	89.	201.0	9.0		2.09	6			
	#	18.	229.7	1.25	-0.003	.	7	6.56		
	7	=	(230)							
РЪО	×	0		. 7		4.37	4.37			ⅎ
	A 200	88	451.7	3,33		1.89	4.38			
19										
	225	က	493.5	2.26		3,33	6.02			
	241	0	(1837)	(3.8)		(2.27)	(2.26)			
	1 257	20	(475)							
	19 BE	Ø	(530) 5	(2,92)		(5,89)	(6.77)			
	346	œ	(044)						>4.86	59
PbS	×	0	428.14	.20		•	. 7			9
	187	68,	261,09	.36		•	۲.			
	217	7.4	282.17	.85		•	S			
	231	50.7	303,93	1.436		2.00	4.87			
	249	52.	283.95	.17		•	. 2			
•	9.5	86.	299.34	. 57		•	•			
									4.55	2 9
	477	33		(7.8)						

TABLE II. (Continued)

	ā						Convergence Limit	Limit	
Molecule	state voo	3	x X	γe	ω ² /4×ω	linea	graph	observ.	Ref.
ı	cm-1	T HO	cm_t	CIII I	ev	ev	ev	ev	
PhSe	0 ×	277.6	0.51		69 ° tı	69 1			59,61
i i	18661.	166,9	0.14		6,03	8,35			•
	В 20959,4	184.8	0.43		2,46	5.06			
	23268.	183.0	0.25		4.15	7.04		,,	
	28374.	190.4	0.53	±00.0-	2.12	5.64	(4.75)	G(
	F 45194.5	224.8	0.50		3.13	8.73		•	
PbTe	o ×	211,96	0.43		3.17	0.17			59,62
_		127,08	0.105		4.77	7.05			•
	19704.	144.	9440		1.46	3.90			
	27141.	142.	1,58		0,40	3.76			
		176.	1.0		0.92	60.9			
	46515.	159.	1.4		09.0	6.37			

a. In most cases the letter label is that of the ref. quoted; it may be different from that used in ref.58. b. Numerically calculated using the formula $E=\omega v+\kappa\omega v^2+y\omega v^3$.

20

This comparison will be started with the E state, since this is the only state for which in general accurate values for the convergence limit are available.

It had already been concluded, on the basis of the agreement between the convergence limits for this state and the Birge Sponer value for the dissociation energy of the ground state, that this state correlates with $Me(^3P) + X(^3P)^{(3)}$.

Inspection of the multiplet splittings given in Table III, shows that, in order to decide on the basis of a comparison of spectroscopic and thermochemical data, which Me(³P_{0,1,2}) and X(³P_{2,1,0}) sublevel are involved the more favorable cases would be CSe, CTe, SiSe, SiTe and perhaps GeSe and GeTe for X(³P) and SnO, SnS, PbO and PbS for Me(³P). Unfortunately, the available data preclude such comparison for CSe, CTe, SiSe, SiTe and GeSe. This comparison can however be made for GeTe, SnO, SnS, PbO, PbS and also SnSe, SnTe and to a certain extent PbSe.

It has already been shown for SnS and PbS that such a comparison leads to the conclusion that the E state correlates with $Me(^3P_1) + X(^3P_{1.0})^{(6)}$.

TABLE III. Excitation energy of ${}^3P_{0,1,2}$ and 1D_2 states of the elements of Group IV and VI(in ev).

	³ P ₀	³ P ₀	³ P ₂	¹ D ₂
С	0	0.002	0.005	1.27
Si	0	0.01	0.03	0.78
Ge	0	0.07	0.18	0.89
Sn	0	0.21	0.43	1.07
РЪ	0	0.97	1.32	2.68
0	0.03	0.02	0	1.97
S	0.07	0.05	0	1.15
Se	0.31	0.25	0	1.19
Te	0.58	0.58	0	1.31

This conclusion is confirmed by a similar comparison for the molecules SnO, SnSe and PbO and to a certain extent GeTe, SnTe and PbSe, as shown in Table IV.

The only molecule for which there seems to be a discrepancy between the dissociation energy so derived from the spectroscopic data of the E state and the thermochemical value is SnTe. It may however be noted that the spectroscopic data for the E state of all tellurides are not satisfactory. For SnTe only a region of apparently continuous absorption is known and interpreted as a possible convergence limit of the E state leading to the value 4.07 e.V. which according to the authors (57) "should be treated with considerable reserve".

The accuracy of the thermochemical data makes it difficult to distinguish between the $X(^3P_1)$ and $X(^3P_1)$ sublevels, the separation of which is maximum 0.06 e.V. for X = Se. From the retational analysis of several states of the molecule PbO, Barrow, Deutsch and Travis (4) have concluded on the basis of the Ω correlation rule and the non-crossing rule, that the most likely correlations for several 0^+ states of this molecule are as indicated in Table IV. Professor Barrow has kindly informed us that a similar analysis for SnO also gives as most likely correlation for the E state, $Sn(^3P_1) + O(^3P_1)$.

The interpretation of the rotational analysis and the comparison of the thermochemical and spect*oscopic data are thus in agreement and in fact confirm and supplement one another.

The conclusion that the E state of the group IV-group VI molecules correlate with $Me(^3P_1) + X(^3P_1)$ has therefore been extended in Table V to the molecules SiS, SiSe, SiTe GeS and GeSe for which accurate convergence limits of the E state and no thermochemical data are available. The dissociation energies of SiS and GeSe so obtained were already used in the section dealing with thermochemical data to derive the heats of formation of the corresponding compounds.

If the correlation of the E state thus seems established it is less easy to derive the correlations of the other states. The fact that most combinations of $Me(^3P_1) + X(^3P_1)$ sublevels can lead to one of several $\Omega = 1$ molecular states make it difficult to propose off hand a correlation for the B(1) and D(1) states. The convergence limit of the B state for SnTe, 3.34 e.V., calculated from the formula $E = \omega v + x\omega v^2 + y\omega v^3$ with $\omega = 230.3$, $x\omega = 1.53$, $y\omega = -0.013$ cm⁻¹ would indicate that this state correlates with $Sn(^3P_0) + Te(^3P_2)$. The B state in PbTe, known up to 2.65 e.V. can however not correlate with $Pb(^3P_0) + Te(^3P_2)$ at 2.23*0.10 e.V. and should at least correlate with $Pb(^3P_0) + Te(^3P_2)$ or any combination

 $Pb(^3P_1)$ + $Te(^3P_2,1,0)$. If the convergence limit of SnTe is considered to be low, but not off by more than 0.5 e.V; and if the B states of both molecules is analogous as would be indicated by the same ratio ω_B/ω_X , then the probable correlation of the B state would be $Me(^3P_0)$ + $X(^3P_1)$. (Table IV).

TABLE IV. Correlation of the X, A, B, C, D and E States of the Group IV-Group VI MeX molecules.

Atomic Me	States X	Possible Ω values	likely mole cular states
3 _{P0}	3 _{P2}	0 ⁺ , 1, 2	x(o ⁺)
Ü	3 _P 1	0~, 1	B(1)?
	3 _P 0	0+	A(0 ⁺)
³ P ₁	3 _{P2}	0 ⁻ (2), 0 ⁺ , 1(3), 2(2), 3	c(o ⁺)
•	3 _{P1}	0-, 0+(2), 1(2), 2	E(O ⁺)
	3 P ₀	0-, 1	D(1)?
³ P ₂	3 _{P2}	0-(2), 0+(3), 1(4), 2(3), 3	(2),4
-	3 _P 1	0-(2), 0+, 1(3), 2(2), 3	
	³ P ₀	0 ⁺ , 1, 2	

It is even less easy to speculate about the correlation of the D state. In PbTe, the difference between the convergence limit of this D state, 3.67±0.20 e.V. obtained from a linear Birge Sponer extrapolation and the thermoche-

mical value for the dissociation energy D_0^{O} = 2.30±0.17 is 1.37 ± 0.37 e.V. It may be noted that the uncertainty quoted on the Birge Sponer extrapolation corresponds to an uncertainty as high as ±70%, which still would not affect the conclusion since the linear dissociation energy of the D state itself is only 0.4 e.V. This difference compared to the excitation energies of the different ${}^{3}P_{0,1,2} + {}^{3}P_{2,1,0}$ combinations for Pb and Te, 0, 0.58, 0.58, 0.97, 1.55, 1.55 1.32, 1.90 and 1.90 e.V., indicates that possible correlations are $Pb(^{3}P_{1}) + Te(^{3}P_{1,0})$ or $Pb(^{3}P_{2}) + Te(^{3}P_{2})$. A selection between those two correlations cannot be made on the basis of the highest vibrational level of the D state observed in any of the MeX molecules. A possible argument for the ${}^{3}P_{1}$ + ${}^{3}P_{1}$ correlation could be that the linearly extrapolated convergence limit of the D state lies either above or below the similarly evaluated limit of the E state when the excitation energy $Me(^{3}P_{1}) + X(^{3}P_{1})$ is either larger or smaller than $Me(^3P_2) + X(^3P_2)$ for 8(SiSe, SiTe, GeSe, GeTe, SnO, SnS, SnSe, SnTe) out of 12 molecules of which two (GeO, GeS) have both limits essentially at the same level and two (SiO, SiS) present the inverse situation.

The F states, have already been shown (3) to correlate with $Me(^1D_2) + X(^1B_2)$. They have not been observed for all molecules of this group and will not be discussed here. For SnSe the combination $Sn(^1D_2) + S(^1D_2)$ leads to

one shallow and one repulsive state, from which accurate values of the dissociation energy can be derived (see footnote c, Table V).

For SnTe, higher molecular states G, H and I are known. Perhaps could these states be correlated with $\mathrm{Sn(}^3\mathrm{P}_2\mathrm{)} + \mathrm{Te(}^3\mathrm{P}_{2,1,0}\mathrm{)}.$ The term values voincrease regularly for the A to I states and an abrupt increment for the F state. The same situation occurs for the other molecules of this group as can be seen seen from Table II.

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TABLE V. Comparison of spectroscopic D" and thermochemical DOvalue of the dissociation energies of group IVgroup VI molecules (in e.V.)

Molecule	Convergence limit E state	Excitation energy Me(³ P ₁)+X(³ P ₁)	D"' O	D _O
CO CS CSe CTe			11.111.50.06	7.53±0.1 6.00±0.1 (4.8 ±0.4
SiO SiS SiSe SiTe	6.44±0.06 5.66±0.15 (4.70±0.25)	0.06 0.26 0.59	6.38±0.06 5.40±0.15 (4.11±0.25)	8.32±0.15
Ge0 GeS GeSe	5.79±0.03 5.25±0.0\$	0.12 0.32	5.67±0.03 4.93±0.03	6.79±0.13 5.78±0.15
GeTe	4.49 +0.30 -0.10	0.65	3.84 ^{+0.30} -0.10	4.00±0.15
Sn0	5.68±0.05	0.23	5.45±0.05 5.37(b)	5.53±0.13
SnS	5.06*0.10	0.26	4.80*0.10 4.74(c) 4.83(c)	4.78*0.12
SnSe SnTe	4.54±0.06 (4.07)?	0.46 0.79	4.08 [±] 0.06 (3.28)?	4.11±0.17 3.45±0.17
Pb0 PbS PbSe PbTe	>4.86 4.55*0.20 >4.32 ?	0.99 1.02 1.22	> 3.87 3.53*0.20 > 3.10?	4.00±0.13 3.42±0.12 2.95‡0.22 2.30±0.17

a estimated

a estimated b from absorption continuum at 1926 to 1937 Å, correlated (3) with Sn(\frac{1}{D_2})+O(\frac{3}{P}). It was assumed here that the O(\frac{3}{P_2}) sublevel is involved. If the absorption occured from v"=1, which could be the case as a result of the large interatomic equilibrium distance in the shallow upper state, the value obtained D"=5.47 eV would be in perfect agreement with D"=5.45 eV.

c from shallow q and repulsive r state respectively, assuming both to correlate with Sn(\frac{1}{D_2}) + S(\frac{1}{D_2})

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